

# Multiferroic $\text{PbFe}_{12}\text{O}_{19}$ ceramics

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**Abstract**  $\text{PbFe}_{12}\text{O}_{19}$  (PFO) powders in hexagonal structure have been synthesized by sol-gel process using lead acetate, glycerin and ferric acetylacetonate as the precursor.  $\text{PbFe}_{12}\text{O}_{19}$  ceramics were obtained by sintering the  $\text{PbFe}_{12}\text{O}_{19}$  powders at 1000°C for 1 h. Distorted flaky hexahedron grains are frequently observed in the SEM images of sintered  $\text{PbFe}_{12}\text{O}_{19}$  ceramics. Large spontaneous polarization was observed in  $\text{PbFe}_{12}\text{O}_{19}$  ceramic at room temperature, exhibiting a clear ferroelectric hysteresis loop. The remnant polarization of  $\text{PbFe}_{12}\text{O}_{19}$  ceramic is estimated to be  $P_r \sim 33.5 \mu\text{C}/\text{cm}^2$ . The Fe oxygen octahedron in its perovskite-like hexagonal unit cell as well as the shift off-center of the small Fe cation are proposed to be the origin of polarization in  $\text{PbFe}_{12}\text{O}_{19}$  ceramics. Meanwhile,  $\text{PbFe}_{12}\text{O}_{19}$  ceramics demonstrate strong ferromagnetism at room temperature. Simultaneous occurrence of large ferroelectricity and strong ferromagnetism in  $\text{PbFe}_{12}\text{O}_{19}$  ceramics holds promise for its application in new generation of electronic devices as a practical multiferroic candidate in single phase.

**Keywords** Multiferroics ·  $\text{PbFe}_{12}\text{O}_{19}$  ceramics · Ferroelectric · Polarization

## 1 Introduction

In recent years, there has been increasing interest in multiferroic materials, which have simultaneous two or

more order parameter in the same phase such as ferromagnetic, ferroelectric and/or ferroelastic ordering [1, 2]. Because the couple between magnetic and electric ordering in multiferroic materials leads to electromagnetic effect, they provide wide range of potential applications such as multiple-state memory elements, novel memory media, transducer and new functional sensor [3].

It is clear now that the materials in which ferroelectricity and ferromagnetism coexist are rare [4, 5] and mostly exhibit rather weak ferromagnetism. Because the room-temperature multiferroism is essential to the realization of multiferroic devices that exploit the coupling between ferroelectric and ferromagnetic orders at ambient conditions,  $\text{BiFeO}_3$  together with more recently revealed  $\text{LuFe}_2\text{O}_4$  [6–11] are currently considered to be promising candidates for practical device applications. The perovskite  $\text{BiFeO}_3$  is ferroelectric ( $T_c \sim 1103$  K) and antiferromagnetism ( $T_N \sim 643$  K), exhibiting weak magnetism at room temperature due to a residual moment from a canted spin structure [12], which could somehow prevent its practical application. Therefore preparation of a material in which large ferroelectricity and strong ferromagnetism coexist would be a milestone for modern electronics and functionalized materials [13]. In the past, scientists were trying to look for multiferroic candidates in the traditional perovskite oxides, the typical example was  $\text{BiFeO}_3$ . However, these perovskite oxides exhibit weak ferromagnetism. In contrast, we start to explore the multiferroic candidates in an opposite direction, to look for ferroelectricity in the traditional ferromagnetic oxides, which holds perovskite-like lattice units.  $\text{PbFe}_{12}\text{O}_{19}$  is the one we are looking for. As the electric polarization was found in materials of  $\text{YMnO}_3$  [14],  $\text{ErMnO}_3$  [15],  $\text{LiNbO}_3$  and  $\text{BiFeO}_3$  with hexagonal structure, it opens up a new direction for potential multiferroic candidate of such magnetic materials as  $\text{PbFe}_{12}\text{O}_{19}$  having hexagonal structure.

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M-type hexaferrites denoted as  $\text{PbFe}_{12}\text{O}_{19}$  has attracted a lot of attention because of their excellent magnetic properties and potential application in various fields [16, 17]. Their properties of large magneto crystalline anisotropy, high saturation magnetization and coercivity result in its usage as plastic magnets, recording media, permanent magnets, and components in microwave and high frequency devices [18]. However, the ferroelectric aspect of  $\text{PbFe}_{12}\text{O}_{19}$  has not been reported yet. In this work, we are going to present the large spontaneous polarization of  $\text{PbFe}_{12}\text{O}_{19}$  ceramic in addition to its strong ferromagnetism at room temperature. The simultaneous occurrence of ferroelectricity and ferromagnetism in PFO ceramics will be discussed in detail in this paper.

## 2 Experimental procedure

Nanocrystalline  $\text{PbFe}_{12}\text{O}_{19}$  powders were prepared by polymer precursor method using lead acetate ( $\text{Pb}(\text{CH}_3\text{COO})_2 \cdot 3\text{H}_2\text{O}$ ) and ferric acetylacetonate as starting materials. Typically, (0.7587 g) lead acetate was dissolved in 15 ml glycerin to form a clear solution. After stirring at  $60^\circ\text{C}$  for 1 h, the solution was distilled in a rotary evaporator at  $120^\circ\text{C}$  for 1 h. In this way, the water in the solution was removed away. The distilled solution was stored in a three neck glass bottle. In order to prevent the hydrolysis of ferric acetylacetonate, the following experiments were carried out in the glove box. (6.3574 g) ferric acetylacetonate was dissolved in 50 ml anhydrous ethanol. Then the storing lead acetate solution was added into the ferric acetylacetonate solution at  $50^\circ\text{C}$  under stirring for 1 h. Here, the molar ratio of lead to iron ions was set to be less than 1:12 to balance the Pb loss. 10 mL ammonia and 15 mL polyethylene glycol mixture solution was added into the above solution, thus a colloid dispersion solution was formed. The dispersion solution was maintained at 353 K under stirring for 8 h. The water and organic solution were removed by centrifugation, the remaining colloid powders were calcined at  $450^\circ\text{C}$  for 1.5 h to completely remove the organic part. Then the powders were calcined again at  $800^\circ\text{C}$  for 1 h to yield pure PFO powder. 0.10 g PFO powders were weighted and pressed in a module into pellet samples, which were then sintered at  $1000^\circ\text{C}$  into PFO ceramics. Phase identification was performed by x-ray powder diffraction (XRD) method with  $\text{Cu K}\alpha$  radiation. The morphology and microstructure of pellet ceramics sintered at various temperatures were determined by SEM images, which are collected by a S-4800 field emission scanning electron microscopy (FESEM). The magnetization was measured using a Quantum Design physical property measurement system (PPMS). For P-E hysteresis loop measurement both surfaces of samples were coated with silver paste as electrodes, the

ferroelectric hysteresis loop was measured using a home-made instrument, termed as ZT-IA ferroelectric measurement system.

## 3 Results and discussion

### 3.1 Structure and microstructure of PFO ceramics

The XRD patterns of the PFO powders being sintered at  $1000^\circ\text{C}$ , together with its standard XRD pattern (PDF#41-1373) are shown in Fig. 1 (a) & (b). The XRD pattern of PFO powders being sintered at  $1000^\circ\text{C}$  matches the standard one very well. The XRD patterns showed the hexagonal structure of pure  $\text{PbFe}_{12}\text{O}_{19}$  (PFO) ceramics, whose lattice parameter was reported to be  $a=5.8945\text{Å}$  &  $c=23.09\text{Å}$ , respectively in the JSPD card. The XRD pattern of  $\text{PbFe}_{12}\text{O}_{19}$  powders presents the hexagonal structure with no extra reflections and is perfectly indexed to (110), (112), (107), (114), (200), (203) and (2,0,11) crystal plane of hexagonal  $\text{PbFe}_{12}\text{O}_{19}$  (PDF# 41–1373). The relative intensities of (107) and (114) peaks, which correspond to the inclined c-axis orientation, are higher than those of (112) and (200) peaks. These results indicate that the grains of  $\text{PbFe}_{12}\text{O}_{19}$  ceramics are randomly oriented.

Figure 2 shows the field emission scanning electron microscopy (FESEM) images of the PFO ceramic being sintered at  $1000^\circ\text{C}$ . The surface of the ceramic was polished before SEM images were taken. The grain size was estimated to be within the range of 0.5–3  $\mu\text{m}$ . The distorted hexagonal flaky grains are frequently observed in the PFO ceramics, as being shown in Fig. 2 (b). Cylinder shaped grains can also be found in the SEM images. The

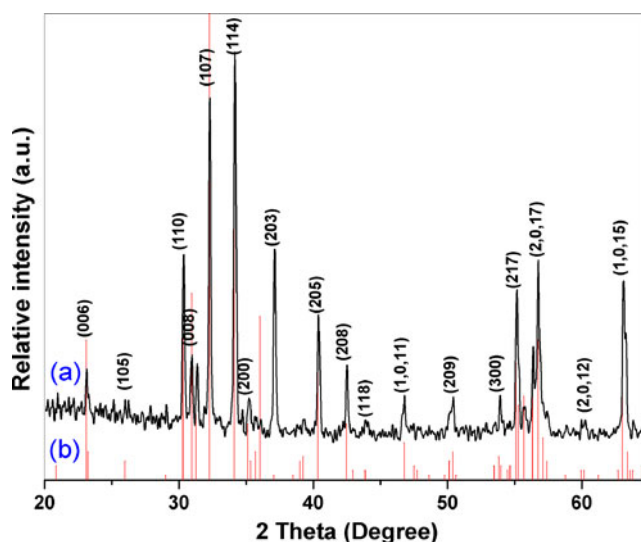
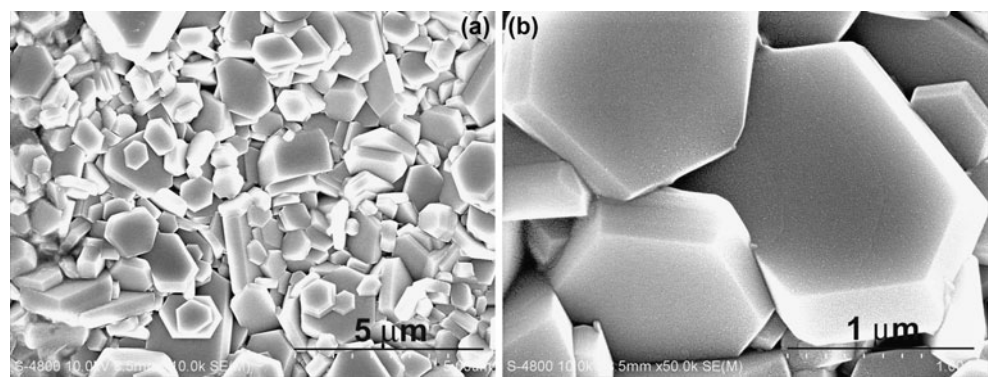


Fig. 1 XRD patterns of the powders calcined at (a)  $1000^\circ\text{C}$  and (b) standard pattern

**Fig. 2** FESEM images of the PFO Ceramic sintered at 1000°C



hexahedron shaped grains are consistent with the hexagonal symmetry of PFO crystal. [0001] direction is normal to the surface of the flaky grains. The flaky shape suggests that the grains didn't take priority growth along with [0001] direction but with [0100] direction, which is the longest side of the flaky grains. Obviously the growth rate of the grains along with [0100] direction is much faster than that along with other directions. The deformation of the grains away from normal hexahedron reflects the distortion of hexagonal unit cells in PFO crystals. None centralized symmetry could exist in such distorted hexahedron grains.

### 3.2 Ferroelectric properties of PFO ceramics

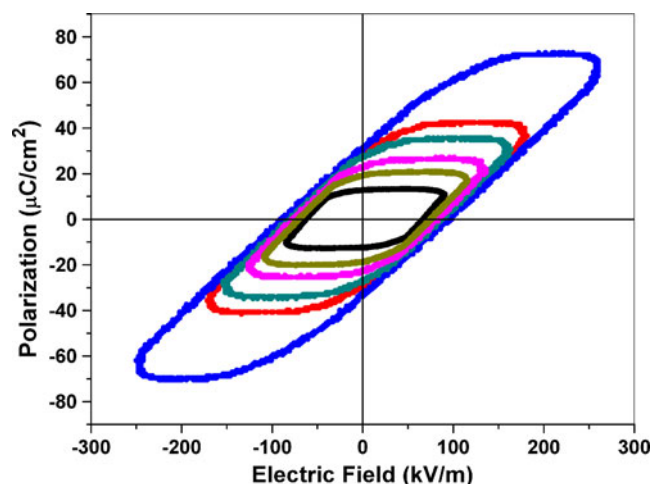
We investigated the effect of the constrained crystallographic state in the PFO ceramics sintered at 1000°C on its physical property. Ferroelectric properties were characterized using polarization hysteresis and pulse polarization measurements. The specimen was parallel connected with a capacitor of 0.1  $\mu\text{F}$  for compensation when ferroelectric measurement was carried out at a frequency of 120 Hz for the PFO ceramics being sintered at 1000°C. Evidence for the characterization of ferroelectric state of PFO ceramic is provided in Fig. 3, which shows polarization cycles exhibiting clear ferroelectric hysteresis loops under applied electric fields of different amplitudes obtained at room temperature. The  $\text{PbFe}_{12}\text{O}_{19}$  ceramics exhibited a quickly increase of saturated polarization with the increase of applied electrical field. It can be seen from the electric hysteresis loops that the maximum remnant polarization ( $P_r$ ) is determined to be  $\sim 33.5 \mu\text{C}/\text{cm}^2$ , which is around 5.5 times higher than the reported value of  $6.1 \mu\text{C}/\text{cm}^2$  from  $\text{BiFeO}_3$  ceramics [19]. The polarization maximum ( $P_{\text{max}}$ ) and the coercive electric field ( $E_c$ ) obtained from the ferroelectric hysteresis loop are  $\sim 70 \mu\text{C}/\text{cm}^2$  and  $\sim 96 \text{ kV}/\text{m}$ , respectively.

To confirm the ferroelectricity of the  $\text{PbFe}_{12}\text{O}_{19}$  ceramics, the polarization characteristic was measured under a pulsed trapezoidal probe condition, which is less likely to be convoluted by leakage and nonlinear dielectric effects. The pulsed remnant polarization (defined as

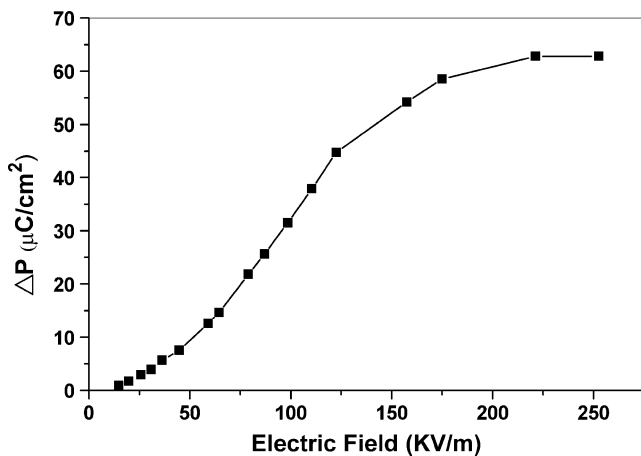
$\Delta P = P^* - P^\wedge \approx 2P_r$ , where  $P^*$  is the switched polarization and  $P^\wedge$  is the nonswitched polarization) shows a sharp increase of  $\Delta P$  before 50 kilovolts, reaching a saturation value of about  $65 \mu\text{C}/\text{cm}^2$  at 221  $\text{KV}/\text{m}$ , as being shown in Fig. 4. This polar state was found to be replicable, as evidenced by repeated experiments over several weeks.

Another evidence to verify the ferroelectric characteristics of PFO ceramics is the polarization measurement at different frequencies. The value of the polarization at zero fields was not scaling with the frequency, which means that it does correspond to a remnant polarization rather than to a polarization due to leakage. Secondly, the polarization measurement at different frequencies within the range of 33~120 Hz and a fixed electric field (100 V) demonstrates the reduction of remnant polarization with increasing frequency, which further confirms the ferroelectrics of PFO ceramics.

The ferroelectrics usually exists in perovskite-structure oxides, such as  $\text{SrTiO}_3$ ,  $\text{PbTiO}_3$ , which have the  $\text{BO}_6$  oxygen octahedron. The ferroelectrics are also found in hexagonal compounds like  $\text{LiNbO}_3$ , which has two  $\text{NbO}_6$  oxygen octahedrons in a unit cell. Below the Curie



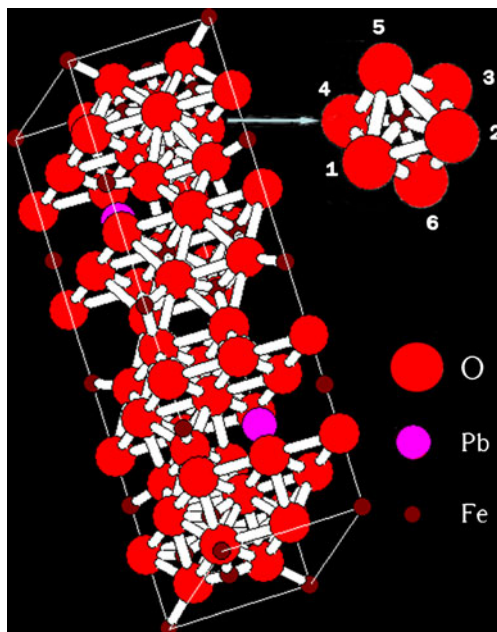
**Fig. 3** Ferroelectric hysteresis loops measured at a frequency of 120 Hz, show that the ceramics sintered at 1000°C is ferroelectric with  $P_r \sim 33.5 \mu\text{C}/\text{cm}^2$



**Fig. 4** P olarization  $\Delta P$  versus electric field measured with electrical pulses of 120 Hz

temperature, there is a distortion to a lower-symmetry phase accompanied by the shift off-center of the Nb cation and the shift off-public face of Li cation along the c-axis.

We firstly investigated the crystal structure model of  $\text{PbFe}_{12}\text{O}_{19}$  (ICSD #36259) with space group  $P63/mmc$ , which is exhibited in Fig. 5. Careful analysis of the unit model structure suggests a perovskite-like crystal structure with one distorted  $\text{FeO}_6$  octahedron in hexagonal  $\text{PbFe}_{12}\text{O}_{19}$  as being shown in Fig. 5. Each hexagonal  $\text{PbFe}_{12}\text{O}_{19}$  model has one  $\text{FeO}_6$  octahedron in a sub-unit cell. In a normal octahedron, Fe cation is located at the center of an octahedron of oxygen anions. However, in the unit cell of  $\text{PbFe}_{12}\text{O}_{19}$  below the Curie temperature, there is



**Fig. 5** Crystal structure model of  $\text{PbFe}_{12}\text{O}_{19}$  (ICSD #36259) and magnified oxygen octahedron perovskite structure found in the crystal structure model of  $\text{PbFe}_{12}\text{O}_{19}$

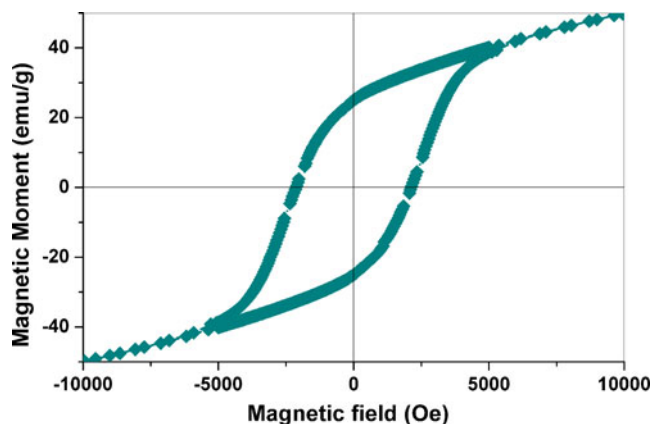
also a distortion to a lower-symmetry phase accompanied by the shift off-center of the small Fe cation. Fe cation shifts away from the center along b axis, while  $\text{O}_3$  and  $\text{O}_4$  shifts off their original positions of octahedron along opposite directions of a-axis, which leads to the distortion of  $\text{O}_5\text{-Fe-O}_6$  bond away from straight line. The spontaneous polarization derives largely from the electric dipole moment created by the two shifts.

The non-central symmetrical octahedron in the hexagonal unit cell is proposed to be the origin of the electric dipole moments and thus be responsible for the spontaneous polarization of the PFO ceramic to the external electronic field, which behaves in large ferroelectric hysteresis loops in Fig. 3.

### 3.3 Ferromagnetism of PFO ceramics

Now let's turn our attention to magnetic response.  $\text{PbFe}_{12}\text{O}_{19}$  is a traditional ferromagnetic material, whose magnetization behavior has been widely studied. Magnetic properties of the PFO ceramics were measured at room temperature with a Quantum Design physical property measurement system (PPMS). The field-dependent magnetization hysteresis loop for PFO ceramic is shown in Fig. 6. The remnant magnetic ( $M_r$ ) of the PFO ceramic was 24.5 emu/g, the saturation magnetization was  $\sim 50$  emu/g. The magnetic coercivity ( $H_c$ ) of the PFO ceramics is 2168 Oe, which is lower than the reported value of 3800 Oe of the PFO ceramics sintered at  $800^\circ\text{C}$  [16]. The PFO ceramic exhibit reduced magnetization value in comparison with its counterpart being sintered at lower temperature.

Size effect plays an important role on the reduction of the magnetization of PFO ceramics. Usually the so-called critical size determines the properties of magnetic materials [19–21]. When particle size is smaller than this critical value, particles locate within one single domain; otherwise multiple-domain may occur in particles. As particles are



**Fig. 6** Magnetic hysteresis loop of  $\text{PbFe}_{12}\text{O}_{19}$  ceramic sintered at  $1000^\circ\text{C}$

larger than the single domain size for the ceramics being sintered at higher temperature, the domain walls become predominant. The larger is particle size, the more are domain walls. The domain walls could prevent somehow the switching of electron spins with external magnetic field, which results in reduction of magnetization of PFO ceramics when particle size increases with sintering temperature.

So far, BiFeO<sub>3</sub> is the best-known multiferroic material, which also exhibits both ferromagnetic and ferroelectric ordering above room temperature. However, the weak ferromagnetism of BiFeO<sub>3</sub> may prevent its application [22]. In contrast, strong ferromagnetism and large ferroelectricity simultaneously occurred in PFO ceramics above room temperature. The remnant magnetic polarization ( $M_r$ ) and the magnetic coercivity ( $H_c$ ) of BiFeO<sub>3</sub> was reported to be less than 0.1 emu/g and 200 Oe [23], which are much less than that of ~24.5 emu/g and 2168 Oe for PbFe<sub>12</sub>O<sub>19</sub> ceramics. The remnant magnetic polarization of PbFe<sub>12</sub>O<sub>19</sub> ceramics is 245 times higher than that of BiFeO<sub>3</sub> ceramics, while the magnetic coercivity ( $H_c$ ) is around 11 times higher than that of BiFeO<sub>3</sub> ceramics. Meanwhile, as being mentioned above, the maximum remnant polarization ( $P_r$ ) of PbFe<sub>12</sub>O<sub>19</sub> ceramics is determined to be ~33.5  $\mu\text{C}/\text{cm}^2$ , which is around 5.5 times higher than the reported value of 6.1  $\mu\text{C}/\text{cm}^2$  from BiFeO<sub>3</sub> ceramics [19]. Therefore PbFe<sub>12</sub>O<sub>19</sub> ceramics have clear and unique advantage over the best multiferroic system of BiFeO<sub>3</sub> materials. This holds promise for its application in new generation of electronic devices as a practical multiferroic candidate in single phase.

#### 4 Conclusion

In summary, pure lead hexaferrite (PbFe<sub>12</sub>O<sub>19</sub>) powders have been synthesized by a novel polymer precursor method using glycerin as solvent. The powders were pressed into pellets, which were sintered into ceramics at 1000°C for 1 h. The ferroelectric hysteresis loop of the sample shows that the remnant polarization ( $P_r$ ) and the coercive electric field ( $E_c$ ) are ~33.5  $\mu\text{C}/\text{cm}^2$  and 96 kV/m, respectively. The polarization  $\Delta P$  versus electric field shows that the saturation polarization ( $\Delta P$ ) of the sample

could reach as high as ~63  $\mu\text{C}/\text{cm}^2$ . We propose that the source of polarization is the distortion of the Fe oxygen octahedron in the lattice unit of its perovskite-like hexagonal structure. The magnetic hysteresis loop of the sample shows that the remnant magnetic polarization ( $M_r$ ) and the magnetic coercivity ( $H_c$ ) are ~24.5 emu/g and 2168 Oe, respectively.

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